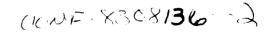
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THE STABILITY OF THE LOW DEGREE FIVE MINUTE SOLAR OSCILLATIONS

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THE STABILITY OF THE LOW DEGREE FIVE MINUTE SOLAR OSCILLATIONS

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In this paper we discuss the decay rate for many of the low degree p modes observed as 5 minute oscillations of the sun. This report is an expanded version of the presentation at Snowmass. These theoretical results use the completely nonadiabatic linear theory of Saio and Cox (1980). Our solar model is based on the evolution results of Christensen-Dalsgaard (1982). Equation of state and opacity data come from the Los Alamos Opacity Library of Huebner, Merts, Magee, and Argo (1977). We compute decay rates for modes ranging from radial (ℓ =0) to the nonradial ones with ℓ =5 for overtones 10 through 28.

Parameters needed for our solar model are given in Table 1. Figure 1 shows the hydrogen mass fraction composition structure. Also given on the figure is the structure given by Christensen-Dalsgaard (1982) for an evolved solar model. Our special equation of state and opacity table with X=0.74 for the hydrogen mass fraction in the outer 0.40 of the mass needs slightly more hydrogen in the central regions than obtained by Christensen-Dalsgaard in order to give a complete and consistent model. The difference in helium production between these two models is about 10%, meaning that the total energy radiated by the sun during its lifetime thus far agrees satisfactorily with accurately calculated evolution sequences.

TABLE 1

SOLAR MODEL:

3.90x10** erg sec-1 Luminosity 1 |89x10** a Mass 6.955x1010 cm Radius 5.8x108 K Surface temperature 1.51×107 K Central temperature 122.3 g cm-Central density 0.740, 0.240, 0.020 Surface X. Y. Z Central X. Y. 2 0.450, 0.630, 0.020 0.32Ro (0.043Mo) Depth of convection zone 2.50x104 K Temp. at bot. of convec. zone 0.05Ro (0.01Mo) Central ball 3.0x10** g Surface-zone mass

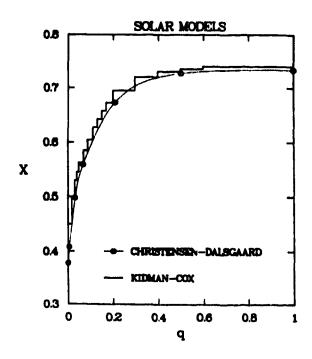


Figure 1. The hydrogen mass fraction of the composition, X, is plotted versus interior mass fraction of the solar model.

Figure 2 is a plot of the logarithms of temperature (K), opacity (cm²/g), and density (g/cm³) for our model versus the logarithm of the exterior mass. The very high opacity over the outer 4% of the mass produces a very deep convection zone. The rapid rise of temperature just cooler than 7,000 K requires a small density inversion to give the proper run of pressure to maintain hydrostatic balance. The ratio of mixing length to pressure scale height for all the convection zone is 1.5.

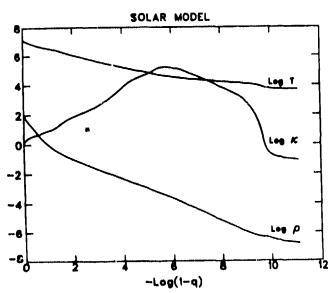


Figure 2. The temperature, opacity, and density structure of the sun is plotted versus the logarithm of the exterior mass fraction.

Solution for the eigenvalues and eigenevectors for the nonradial modes is made for all six of the Dziembowski (1971) variables which have both real and imaginary parts. Figure 3 gives the central variations of y_1 , the Lagrangian variation of the mass shells in the radial direction. The imaginary part, which indicates the variation structure at the mean radius phase of the pulsation as contrasted with the real part applying to the time of the maximum expansion for these linear sinusoidal motions, is very small. This means that the oscillations for this p_{23} , $\ell=2$ mode are very adiabatic. There is little phase change for these lobes which gives essentially standing rather than running waves. Figures 4 and 5 give this same radial variation structure in the outer

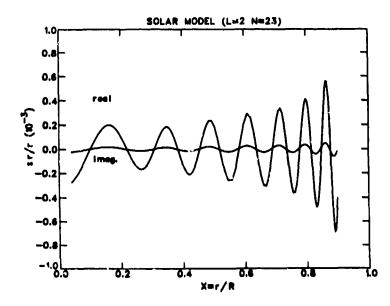


Figure 3. The central variations of the real and imaginary parts of the radial component of the p_{23} &=2 oscillations are plotted versus radius fraction. At x=0.2, 30% of the solar mass is interior. Only 6% of the mass is interior to x=0.1. At x=0.4 the interior mass is 75% of the model mass.

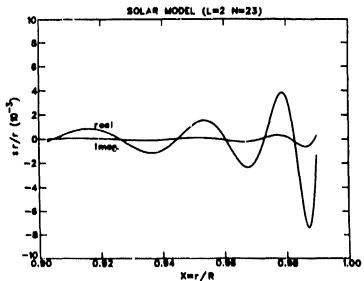


Figure 4. The radial component real and imaginary parts of the p_{23} , $\ell=2$ oscillations are plotted for the outer 10% of the radius.

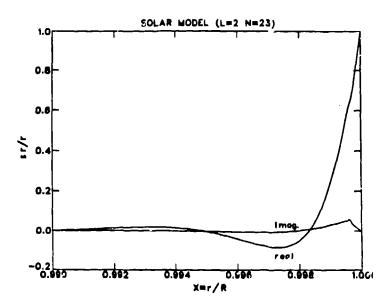


Figure 5. The radial component real and imaginary parts of the p_{23} , $\ell=2$ oscillations are plotted for the outer 1% of the radius.

10% and 1% of the radius. Only at the surface is there any phase change of the real and imaginary parts of this variation, and there also is a significant nonadiabatic effect. At the very surface, the usual normalization gives unity for the real part and zero for the imaginary part. Damping of this mode is related to this feature of the eigensolution, and it is discussed below using the derived complex eigenvalue.

Periods for many modes are given in Tables 2 to 7, respectively, for $\ell=0$ to 5. Modes p_{10} to p_{28} frequencies are listed together with the frequencies observed by Harvey and Duvall (1983). The listed periods and frequencies are from the adiabatic solutions, and the few μ Hz corrections to get the nonadiabatic periods are also given for each mode. Decay rates are calculated by taking the ratio of the imaginary and real parts of the eigenvalue and multiplying this ratio by 4π . It is also possible to integrate the P-V loops that all these Lagrangian zones traverse each cycle to get the work done by or done on these zones each cycle. The sum over all the mass zones agrees well with the decay rate derived directly from the eigenvalue component ratio. For each mode the decay times are given in terms of the number of cycles and in terms of the real time.

Our calculations were done with only 329 mass shells and the central ball. These 330 zones are not quite enough to define well the eigensolutions, and we get frequencies 1.3 to 2.5% too large compared with the observed frequencies. See Figure 5 where a comparison of our frequencies with those of Shibahashi et al. (essentially the observed ones) is made. This trend indicates a zoning error, because it decreases with increasing resolution of the pulsation mode and its eigensolutions. Our optimum zoning has very fine mass shells down to a depth of 3.5 million kelvin. In this outer region the temperature increases only about 3% per zone. In the deeper layers, the 100 more zones have a mass ratio from zone to zone of 1.003. We estimate maybe 100 more carefully placed zones could give adequately accurate periods. Our actual

frequencies should not be accepted for high accuracy, but the decay rates listed in the tables should be well determined.

TABLE 2 periods, frequencies and decay rates

						decay by e	
ı	n	period (sec)	adiabatic frequency (uhz)	non-a incre (uhz)	frequency (uhz)	number of cycles	time (days)
000000000000000000000	11123456789912345678	620.895 573.832 496.392 496.340 438.540 414.175 392.092 372.124 354.083 337.481 322.317 308.399 295.583 283.583 272.244 252.677 243.731	3813.2	-3.39 -3.48	9.0	330.4 269.5 227.6 191.9 167.0 144.7 132.0 120.6 109.4	131.73 556.489 14.291 1.356.281 18.01 1.356.291 1.360

TABLE 3
periods, frequencies and decay rates

			adimbatic	non-a	Avasured	decay	by •
ι	n	period (sec)			frequency (uhz)	number of cycles	time (daya)
Ĭ	1567890123456 1222223456	607.745 600.033 600	1645.4 1782.7 1919.3 2056.8 2193.8 2331.4 2466.8 2602.3 2739.6 2878.0 30158.4 2139.8 3139.8 3139.8 3158.4 4156.9		0.0 0.0 0.0 0.0 2161.0 2293.0 2427.0 0.0 2828.0 2962.0 3098.0 3233.0 3368.0 35641.0 3717.0 4058.0	13691.5 6976.6 73384.7 13994.7 12175.6	95.67 410.422 116.422 116.867 11.8665

TABLE 4
periods, frequencies and decay rates

						decay	by •
ι	n	period (aec)	frequency	non-a incre (uhz)	frequency (uhz)	number of cycles	time (days)
22222	123456789012345678	588.804 544.901 473.968 444.909 419.299 396.417 375.999 357.558 340.615 325.108 310.911 297.790 285.664 274.496 264.189 2545.479 237.019	3928.7 4073.7	-1.54 -2.52 -2.92	3980.0	111.8 107.1 99.4	67.99 16.89 16.89 5.127 1.50 1.152 1.515 1.39 1.32 1.33 1.33 1.33 1.33 1.33 1.33 1.33

TABLE 5 periods, frequencies and decay rates

ι	n	period (sec)	non-a incre (uhs)	measured frequency (uhz)	decay number of cycles	by e time (days)
\star	0123456789012345678	571.766 579.9874 579.9874 579.9874 579.9874 579.9874 579.988 579.9874 579.988 579.98	-1,75 -2.65 -2.39 -2.47	0.0 0.0 0.0 0.0 0.0 0.0 0.0 0.0 0.0 0.0		2088736602448878273075 467136602448878273075 42377421111

TABLE 6
periods, frequencies and decay rates

			adiahakia		measured	decay	by •
ι	n	period (sec)	frequency			number of cycles	time (days)
4	15 16 17 18 19 12 22 23	557.147 517.166 482.1666 451.986 425.986 401.713 380.754 361.671 344.217 328.2627 313.6627 287.614 256.241 246.976 238.338	4195.1	-2.91 -3.29		112.9 106.7 99.9 93.4	37.79 18.71 10.26 53.48 3.28 1.61 1.20 .52 .40 .32 .29 .26

TABLE 7 periods, frequencies and decay rates

ι	n	period (sec)	adiabatic frequency (uhz)	incre	measured frequency (uhs)	decay number of cycles	by e time (daye)
***************************************	1123456789 9 12345678	543.763 545.4871 442.601 442.608 394.088 373.671 3355.377 395.586 272.190 283.587 262.6144 235.287 245.6144 235.287	1833.4 1838.3 1879.3		9.0 1963.0 8100.0 8235.0 82371.0 2605.0 8641.0 8777.0 8914.0 3184.0 3184.0 3187.0 3733.0	4739.6 4739.7 2841.7 9612.0 612.0 612.0 6132.0 64332.0 64332.0 1431.3 1495.3 1495.3 1495.3 1495.3 1495.3 1495.3	29.83 15.32 8.96 8.95 1.97 1.10 2.71 .89 .71 .89 .33 .33 .33 .33 .83

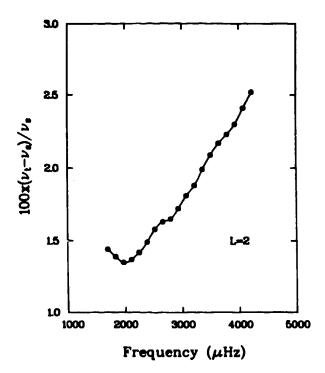


Figure 6. The mode frequency error, as judged from the Shibahashi et al. values, is plotted versus the mode frequency.

A comparison with observations seems to show reasonable agreement with our predictions. The width of the peaks in the power spectrum indicates a decay rate, which is perhaps a matter of days. The coherence of modes over timescales of a month or longer might mean that the mode is reexcited in its existing phase, or it might refer merely to our predicted longer lived modes. Decay rates are faster for higher frequencies because they refer to smaller-scale structures which can more easily gain and lose energy during an oscillation.

Figure 7 shows the work over each pulsation cycle for the outer 30 zones down to a depth of 12,000 K. Actually, the opacity library does not give data for such low temperatures, and the opacities and equation of state are obtained ever this region by use of the Stellingwerf (1975ab) analytic fit. The photospheric damping, always assuming a radiation diffusion structure, down to depths of about 40 Rosseland mean opacity mean free paths is the main damping of the solar oscillations. Zone 317 is at $\tau=10$, zone 319 is at $\tau=5$ and zone 324 is at $\tau=1$. At least half of the damping occurs where our diffusion approximation is valid. The γ and κ effects of the hydrogen ionization region give some deeper pulsation driving, but it is not enough to self-excite the solar oscillations. The convection zone is neutral because we assume, as is usually done, that the convective flux is frozen at its equilibrium value. Deeper than 10,000 K the very small radiative flux is still modulated, but no more hydrogen or helium driving is significant.

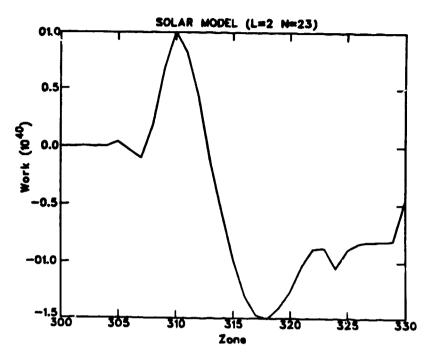


Figure 7. Damping and driving regions of the surface layers are plotted for the outermost 30 zones.

Our conclusions are that all solar 5-minute modes of low order are damped with widely varying rates which increase with frequency. These nonadiabatic effects are confined to the outer 5×10^{-10} of the mass and the outer 3×10^{-4} of the radius.

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